

Particle acceleration in AGN jets

Jieshuang Wang (jswang@mpi-hd.mpg.de)
Max-Planck-Institut für Kernphysik, Germany

In collaboration with Brian Reville (MPIK), Frank Rieger (Uni-Heidelberg),
Yosuke Mizuno (TDLI), Felix Aharonian (DIAS,MPIK)

June 25 @ University of Hong Kong



MAX-PLANCK-INSTITUT
FÜR KERNPHYSIK

UNTERSTÜTZT VON / SUPPORTED BY



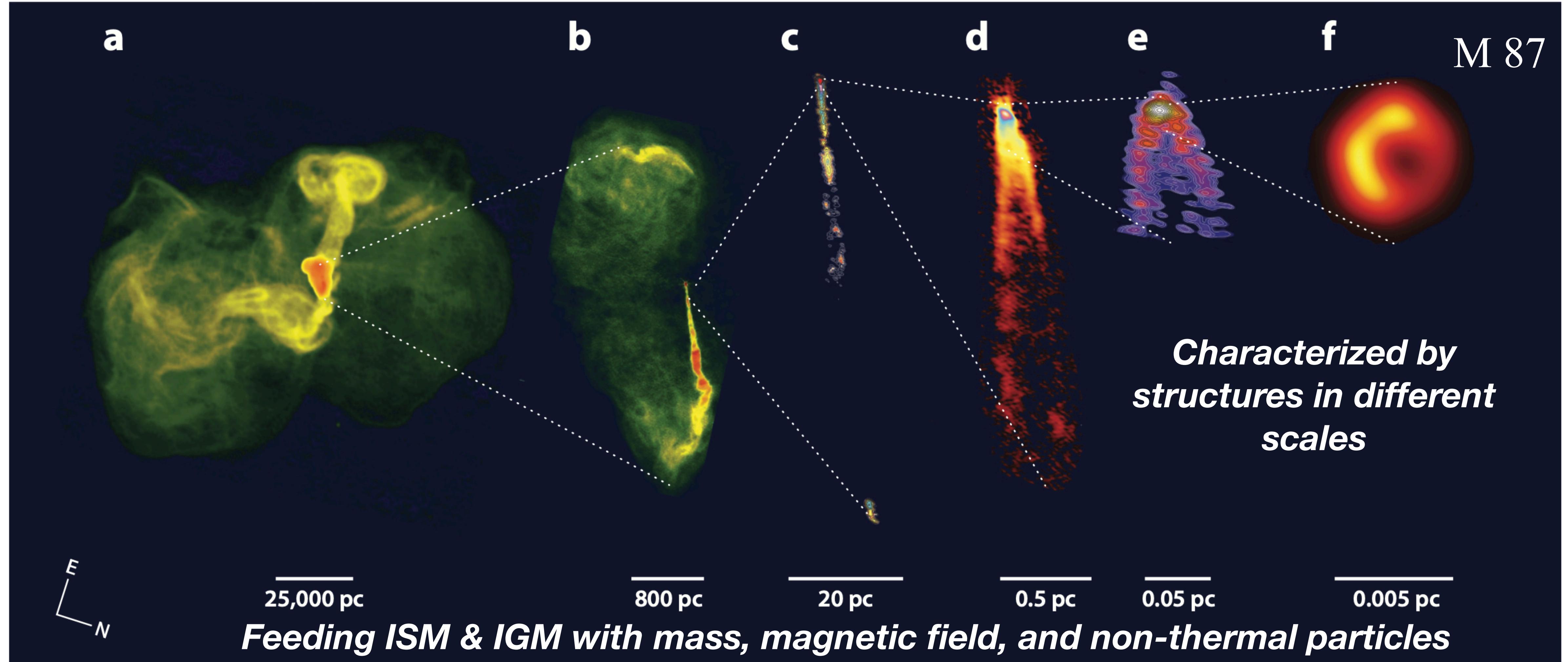
Alexander von
HUMBOLDT
STIFTUNG



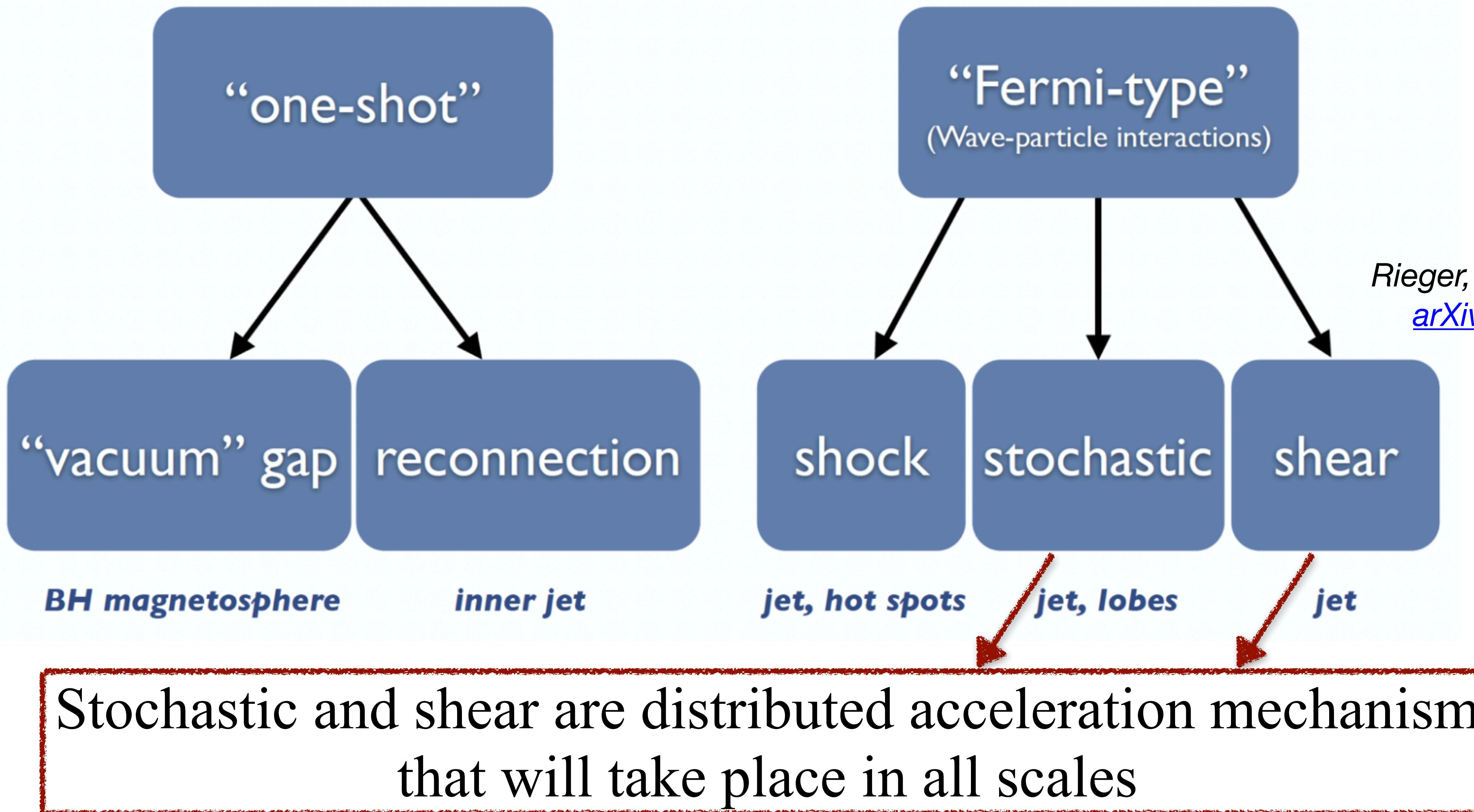
Outline

- Introduction of particle acceleration
- Analytical solutions and applications of shear and stochastic acceleration
- Numerical simulations of particle acceleration
- Summary

AGN jets in all scales



Particle acceleration in different scales



Stochastic and shear acceleration

- Stochastic: standard Fermi-II mechanism
- Energy gain in each collision (Fermi, 1949, Phys. Rev. 75, 578)

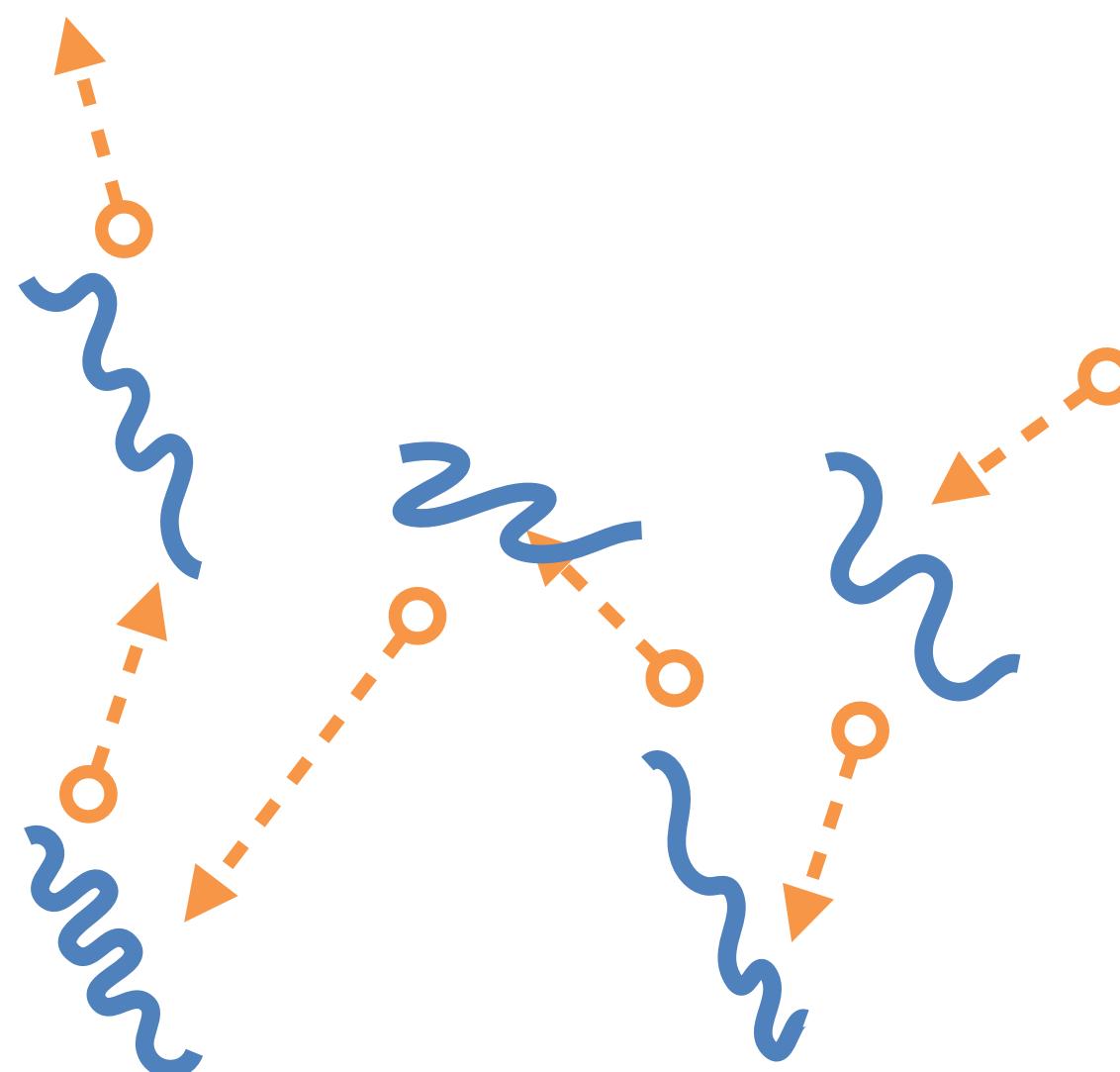
$$\frac{\langle \Delta\epsilon \rangle}{\epsilon} \propto \left(\frac{u}{c}\right)^2$$

- Scattering off MHD waves:

$$u = v_A = B/\sqrt{4\pi\rho}$$

- Scattering time: $\tau_{sc} \propto \gamma^{2-q}$
- Acceleration time (Kolmogorov $q = 5/3$):

$$\tau_{\text{stoch}} = \frac{\epsilon}{\langle \Delta\epsilon \rangle} \tau_{sc} \propto \gamma^{1/3}$$



Stochastic and shear acceleration

- Stochastic: standard Fermi-II mechanism
- Energy gain in each collision (Fermi, 1949, Phys. Rev. 75, 578)

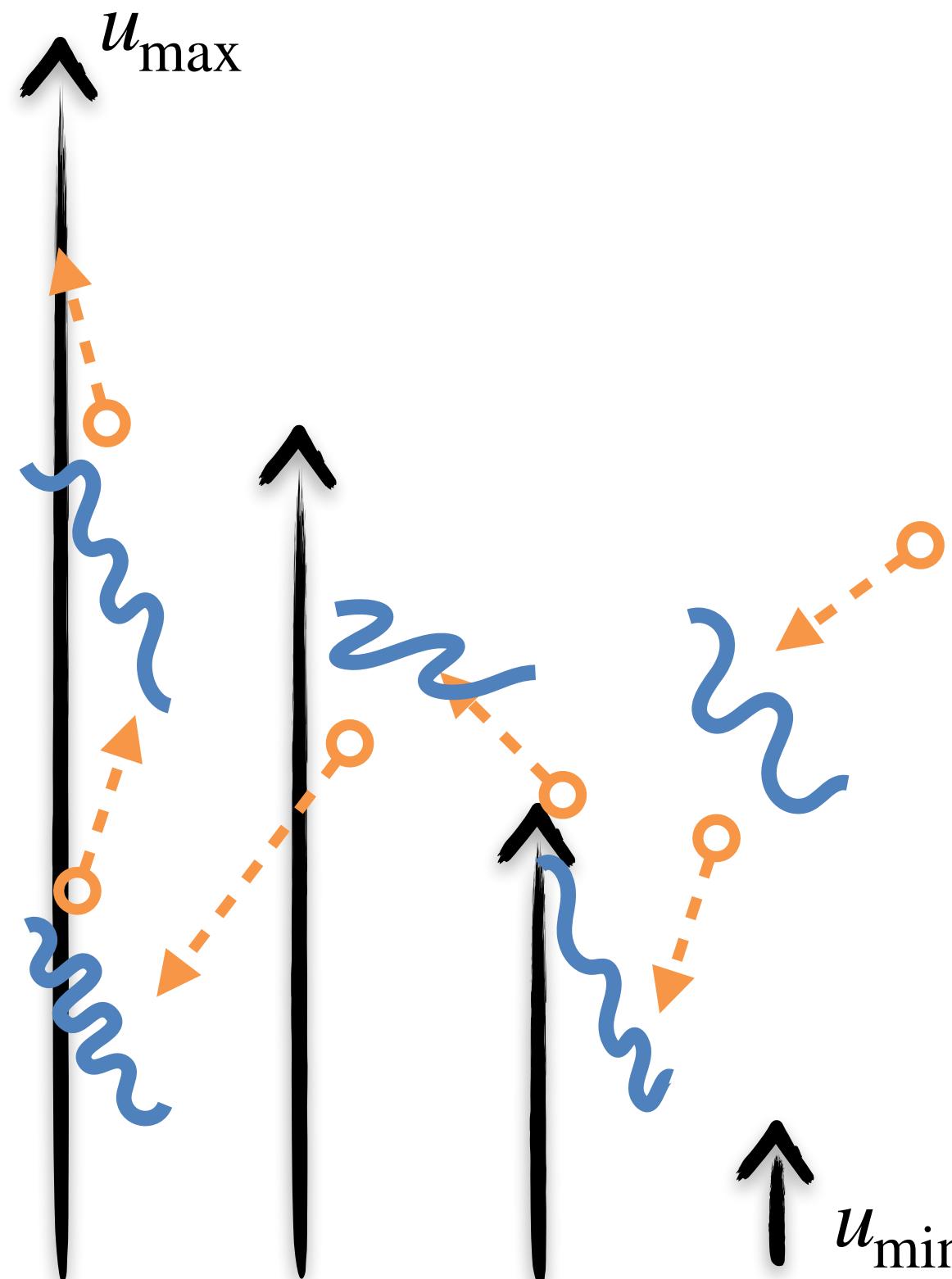
$$\frac{\langle \Delta\epsilon \rangle}{\epsilon} \propto \left(\frac{u}{c}\right)^2$$

- Scattering off MHD waves:

$$u = v_A = B/\sqrt{4\pi\rho}$$

- Scattering time: $\tau_{sc} \propto \gamma^{2-q}$
- Acceleration time (Kolmogorov $q = 5/3$):

$$\tau_{stoch} = \frac{\epsilon}{\langle \Delta\epsilon \rangle} \tau_{sc} \propto \gamma^{1/3}$$



- Shear is also Fermi-II type (see Rieger, 2019, [arXiv:1909.07237](https://arxiv.org/abs/1909.07237) for a review)
- Turbulences are embedded in velocity-shearing layers (spine-sheath)
- Particles scattering off turbulence will sample the velocity difference

$$\frac{\langle \Delta\epsilon \rangle}{\epsilon} \propto \left(\frac{\bar{u}}{c}\right)^2 \propto \left(\frac{\partial u_z}{\partial x}\right)^2 \tau_{sc}^2$$

$$\tau_{shear} = \frac{\epsilon}{\Delta\epsilon} \tau_{sc} \propto \tau_{sc}^{-1} \propto \gamma^{-1/3}$$

Berezhko & Krymsky 1981; Berezhko 1982; Earl+ 1988; Webb 1989; Jokipii & Morfill 1990; Webb+ 1994; Rieger & Duffy 2004, 2006, 2016; Liu+ 2017; Webb+ 2018, 2019; Lemoine 2019; Rieger & Duffy 2019, 2021, 2022



Analytical solutions and application in kpc-scale jets

Particle spectrum from shear acceleration

- An exact solution of Fokker-Planck equation for steady-state shear acceleration:

$$n(\gamma) = C_+ \gamma^{s_+} F_+(\gamma, q) + C_- \gamma^{s_-} F_-(\gamma, q)$$

$$s_{\pm} = \frac{q-1}{2} \pm \sqrt{\frac{(5-q)^2}{4} + w}$$

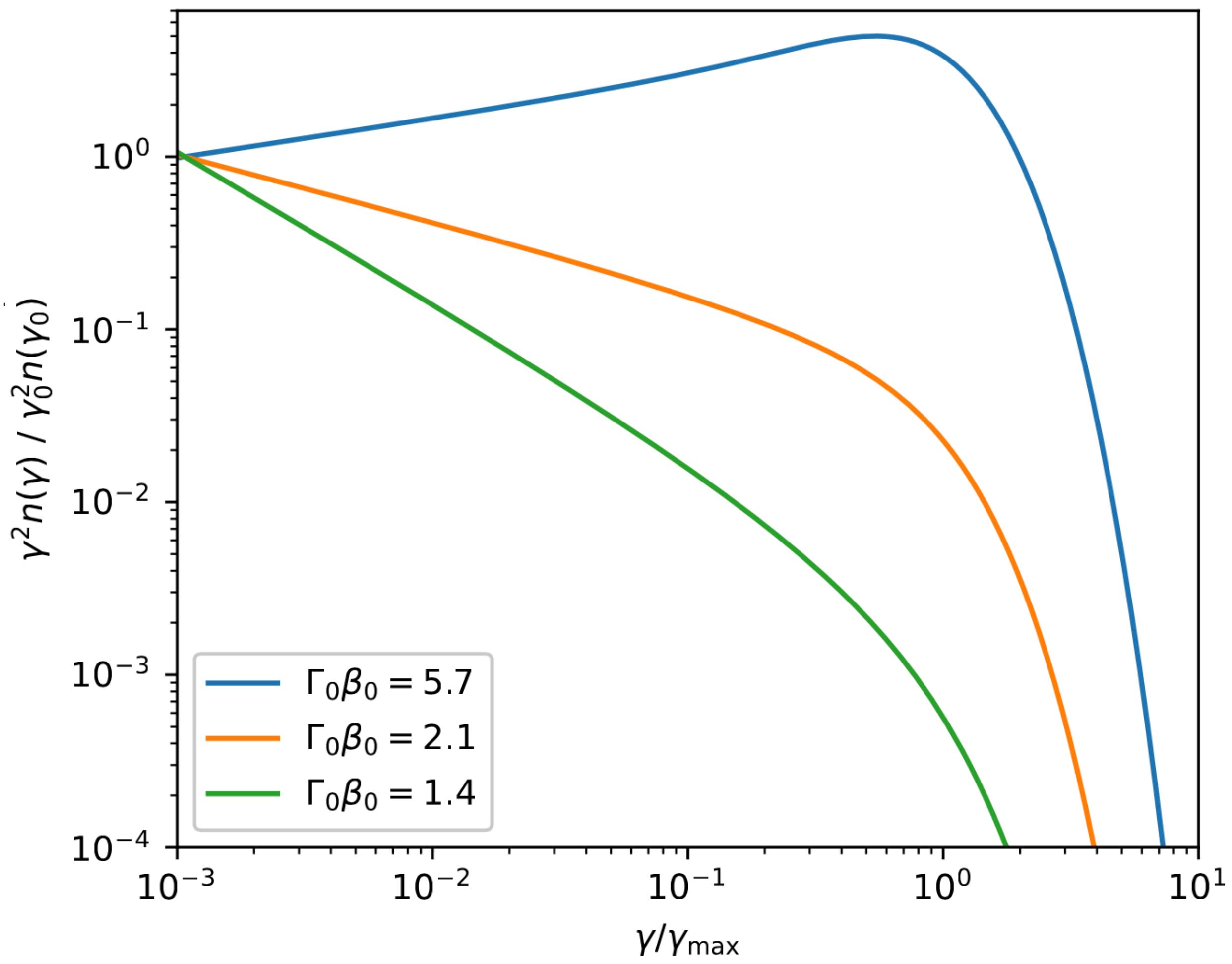
$$F_{\pm}(\gamma, q) = {}_1F_1 \left[\frac{2+s_{\pm}}{q-1}, \frac{2s_{\pm}}{q-1}; -\frac{6-q}{q-1} \left(\frac{\gamma}{\gamma_{\max}} \right)^{q-1} \right]$$

$n \rightarrow 0$ for $\gamma \rightarrow \infty$

- Kolmogorov turbulence: $q=5/3$
- Assume a linear velocity profile

Rieger & Duffy, 2019, *ApJL*, [arXiv:1911.05348](https://arxiv.org/abs/1911.05348)

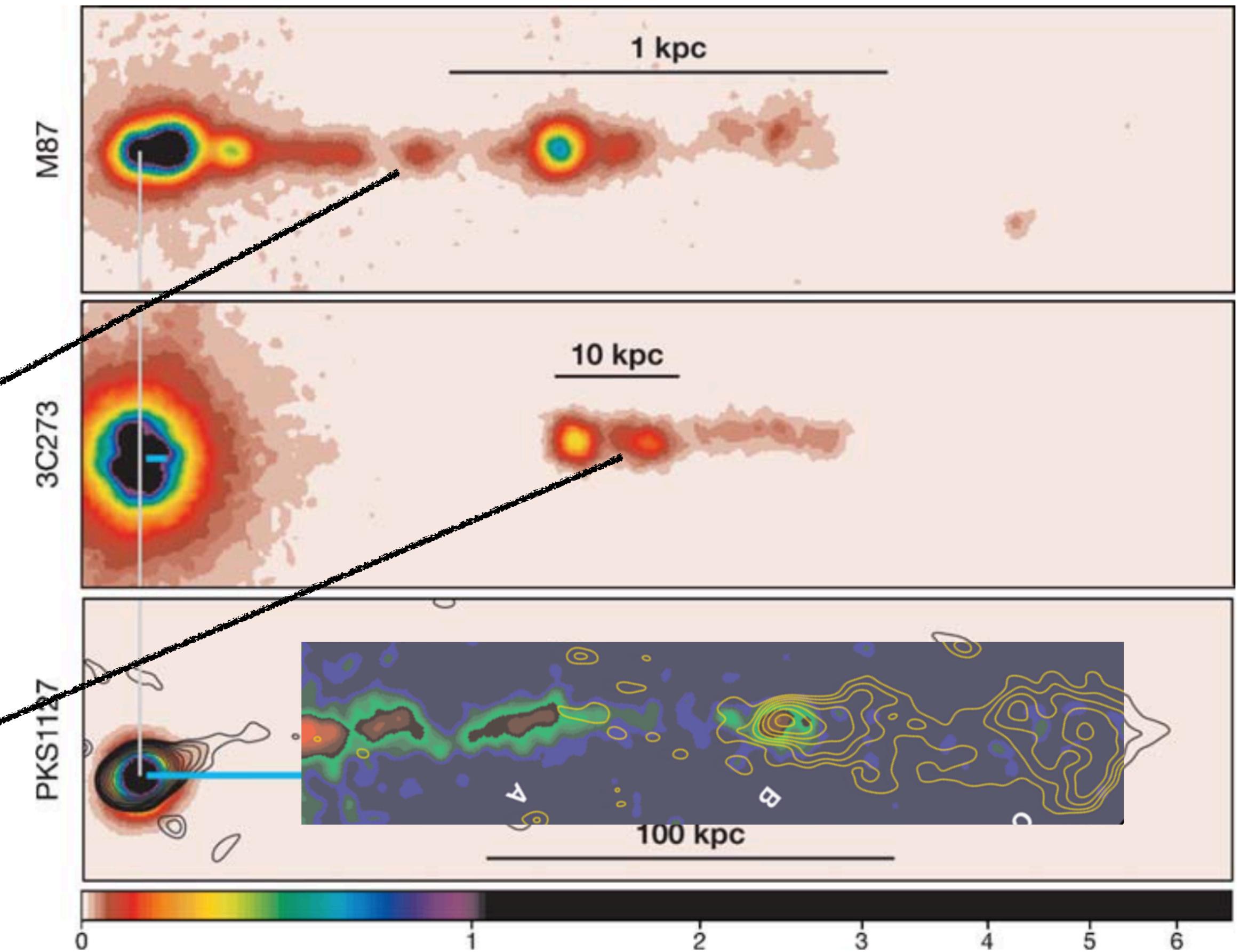
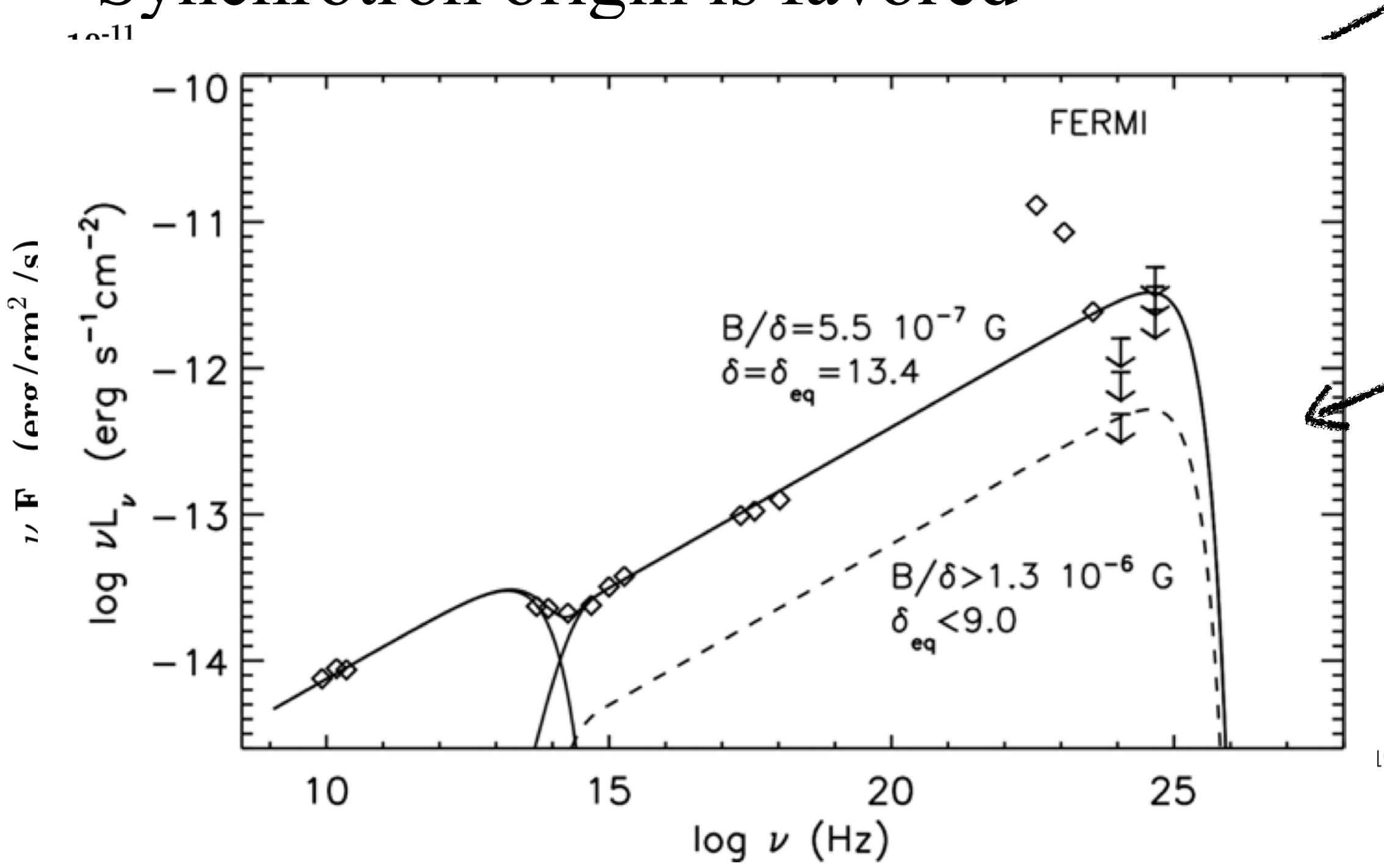
$$w = 40 \ln^{-2} \frac{(1+\beta_0)}{(1-\beta_0)}$$



J.S.Wang+, 2021, *MNRAS*, [arXiv:2105.08600](https://arxiv.org/abs/2105.08600)

kpc-scale X-ray jets

- Quasi continuous X-rays are observed from kpc - 100kpc
- More than 100 sources @ <https://heawww.harvard.edu/XJET/#morph>
- Synchrotron origin is favored



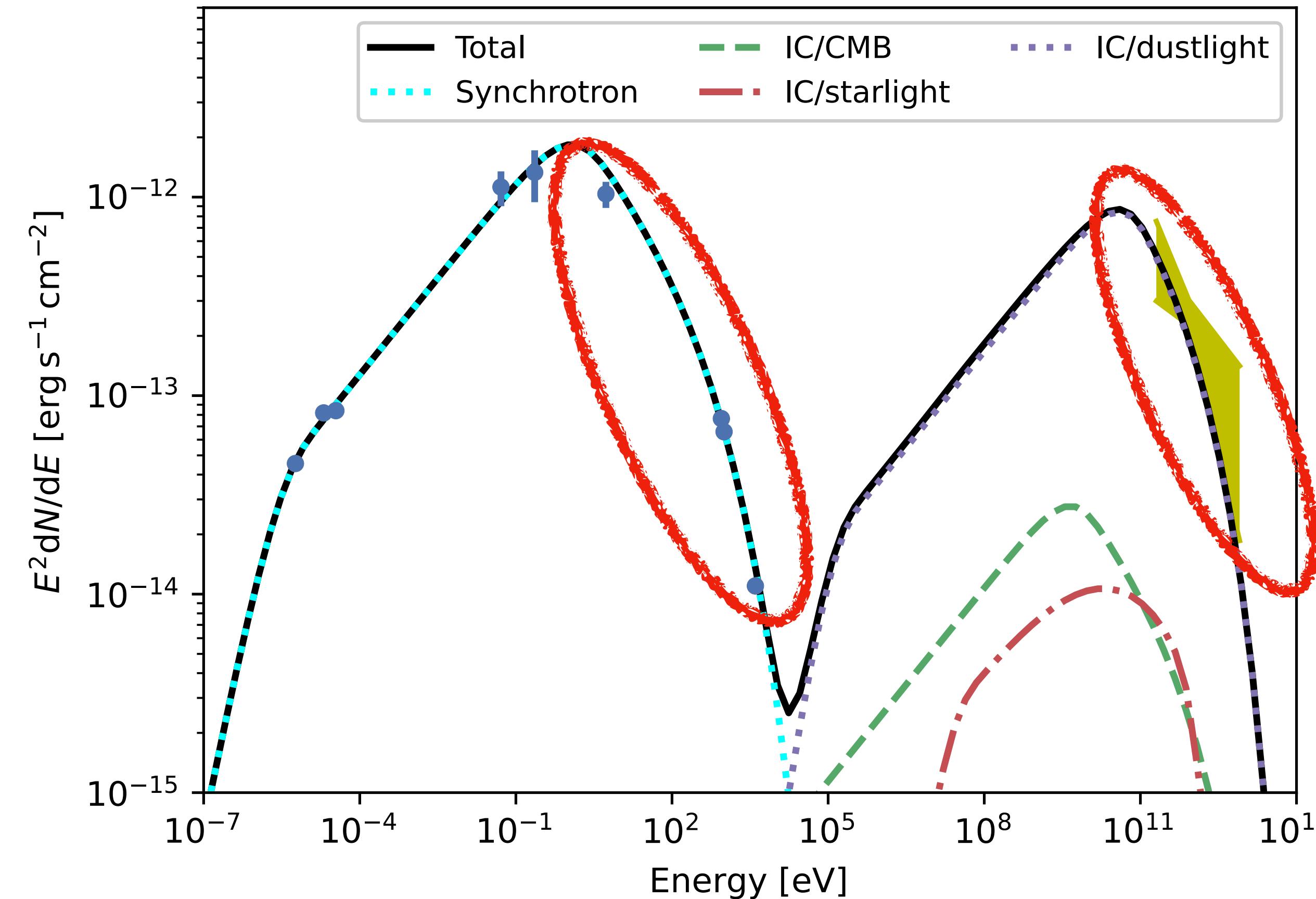
Harris & Krawczynski, 2006, ARA&A, [arXiv:astro-ph/0607228](https://arxiv.org/abs/astro-ph/0607228)

Distributed acceleration required

- Synchrotron origin of X-rays requires sub-PeV electrons ($\text{PeV} = 10^{15}\text{eV}$):
$$E_{\text{syn}} = 2(E_e/0.1\text{PeV})^2(B/10\mu\text{G}) \text{ keV}$$
- Cooling time of sub-PeV electrons:
$$\tau_{\text{syn}} = 1.2 \times 10^3(B/10\mu\text{G})^{-2}(E_e/0.1\text{PeV})^{-1} \text{ yrs} \rightarrow \text{maximum travel distance } c\tau_{\text{syn}} = 0.37 \text{ kpc}$$
- For jet length $>$ kpc, particles accelerated by the jet head shock will cool down immediately after the shock passes (standing shocks may only exist in specific locations)
- *In-situ* (re-)acceleration mechanisms are required along the jet
 - Application of analytical solution of shear acceleration

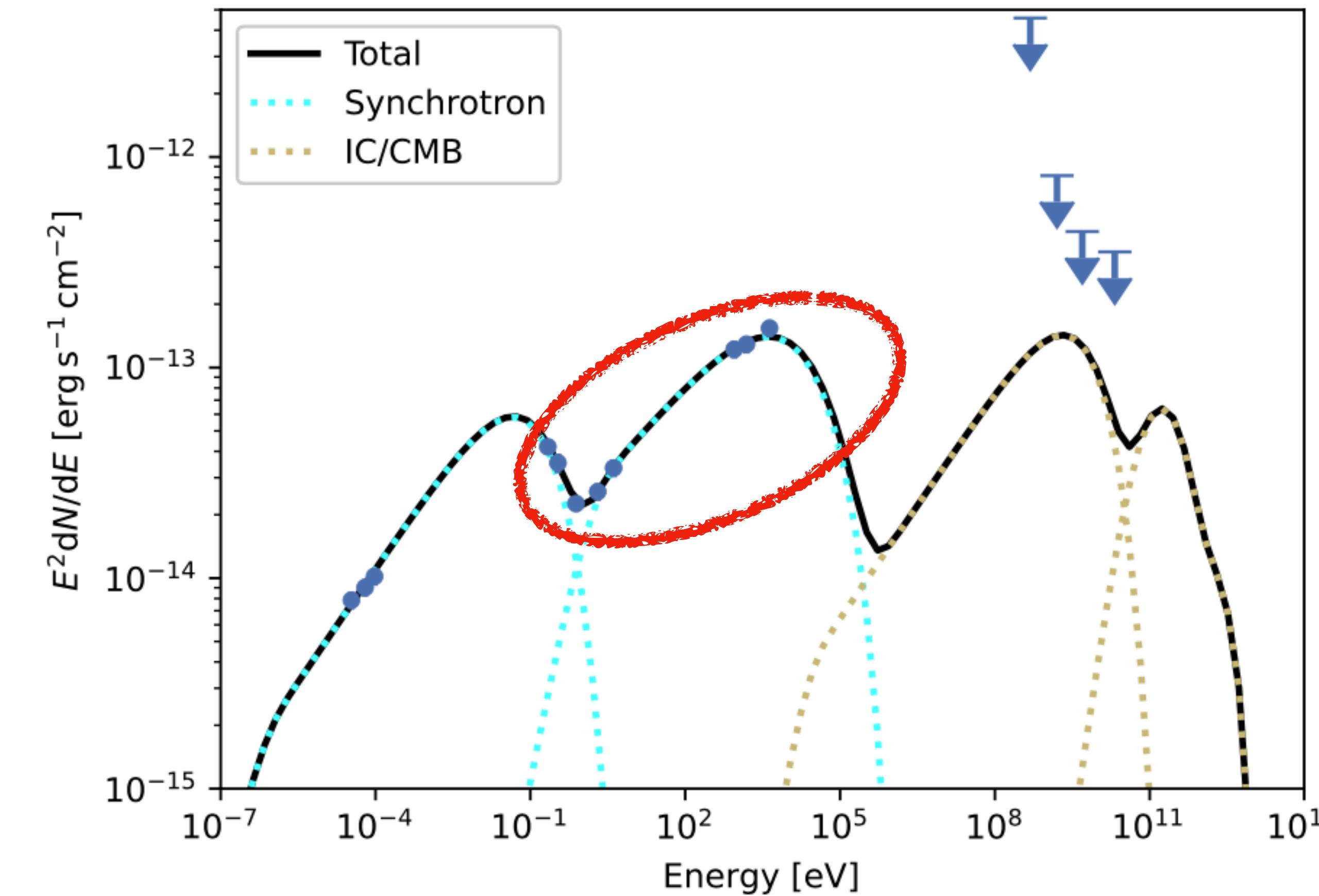
Applications of analytical solution on shear

FR I: Centaurus A



IC and Syn from one population of electrons

FR II: Knots A+B1 of 3C 273

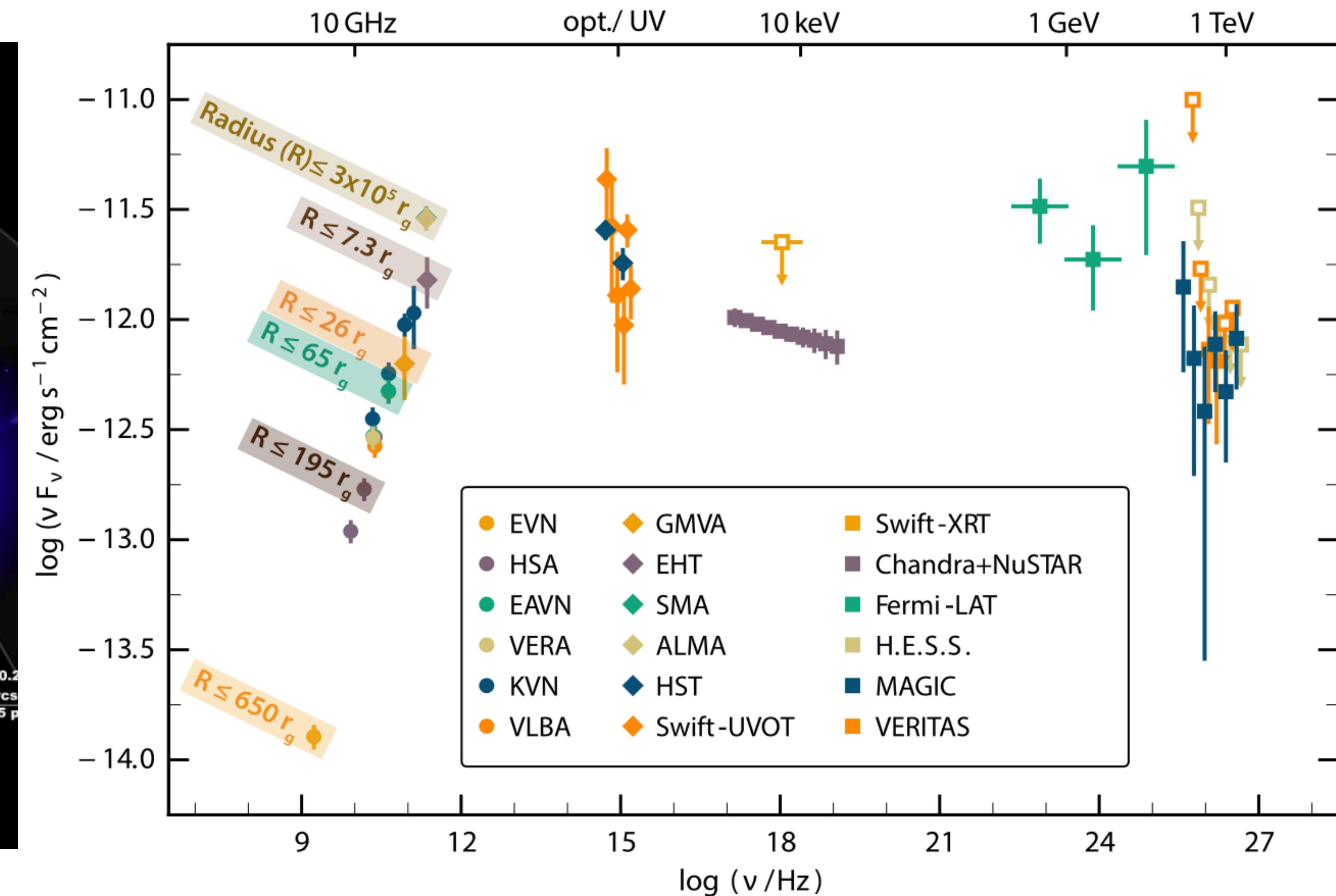
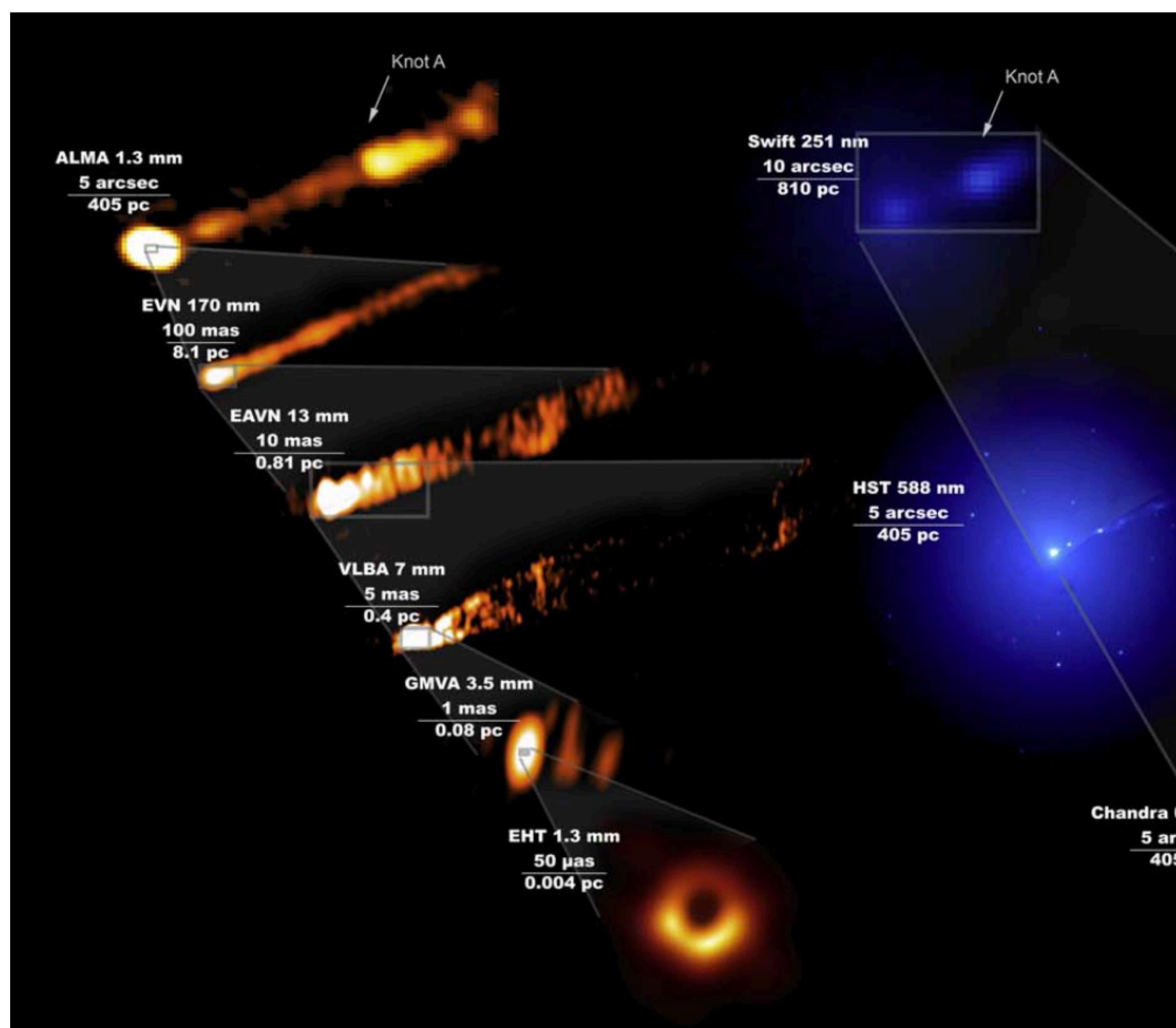


Synchrotron and IC from two populations of electrons

Analytical solutions and application in sub-pc-scale jets

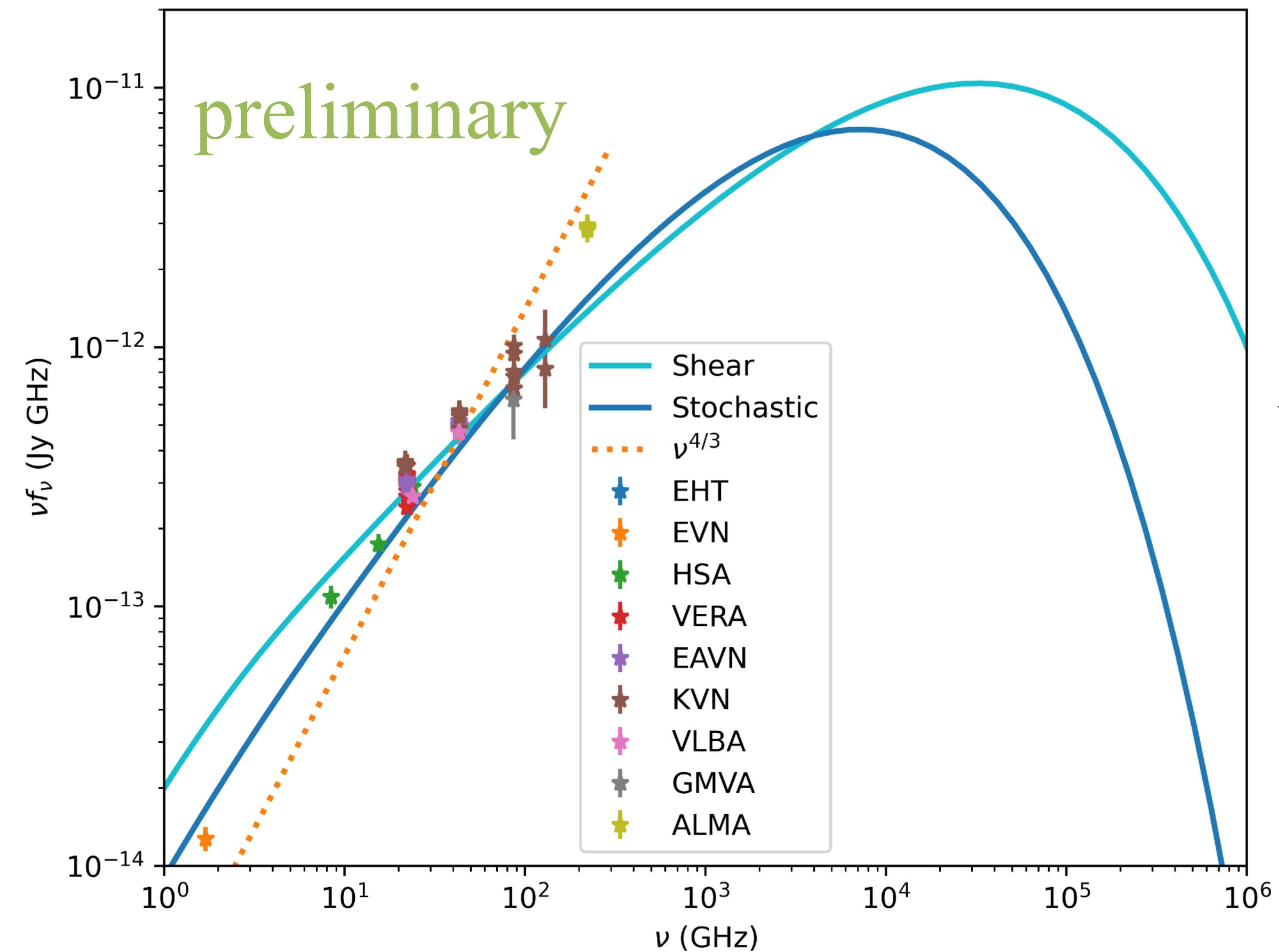


sub-pc-scale radio jet of M87

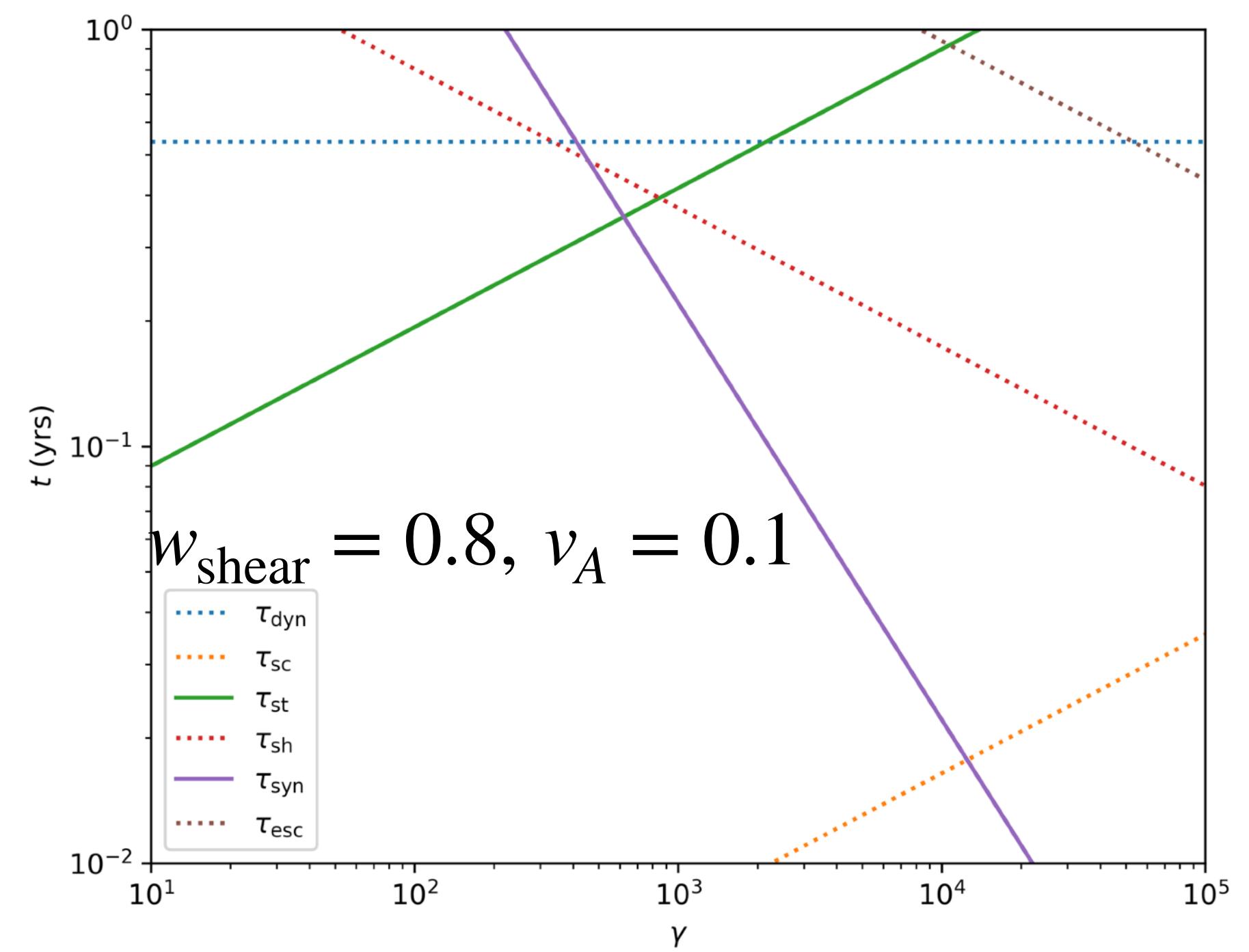
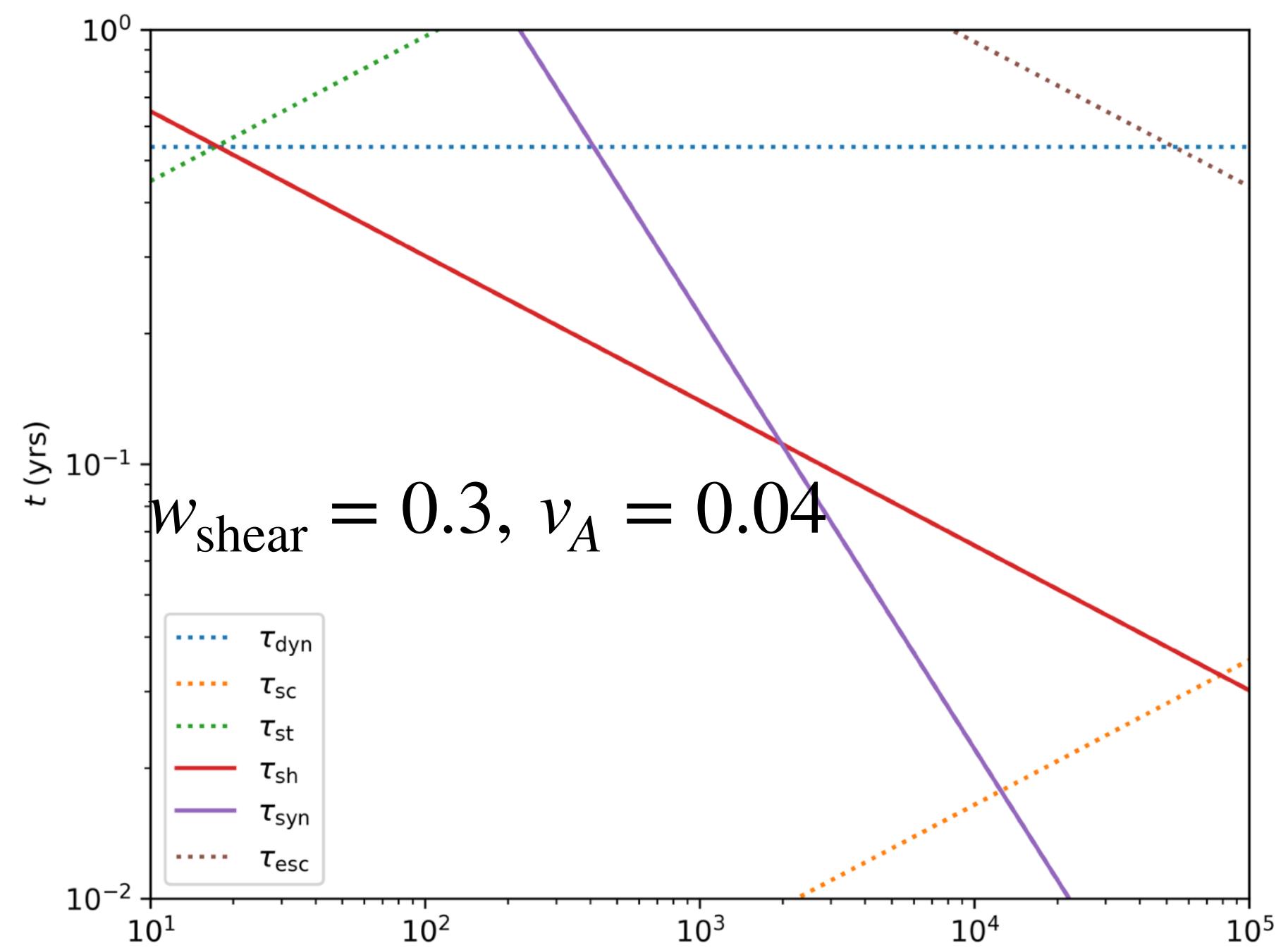


EHT MWL Science Working Group, 2021, ApJL, [arXiv:2104.06855](https://arxiv.org/abs/2104.06855)

Shear and/or stochastic



$$B' = 0.3G$$



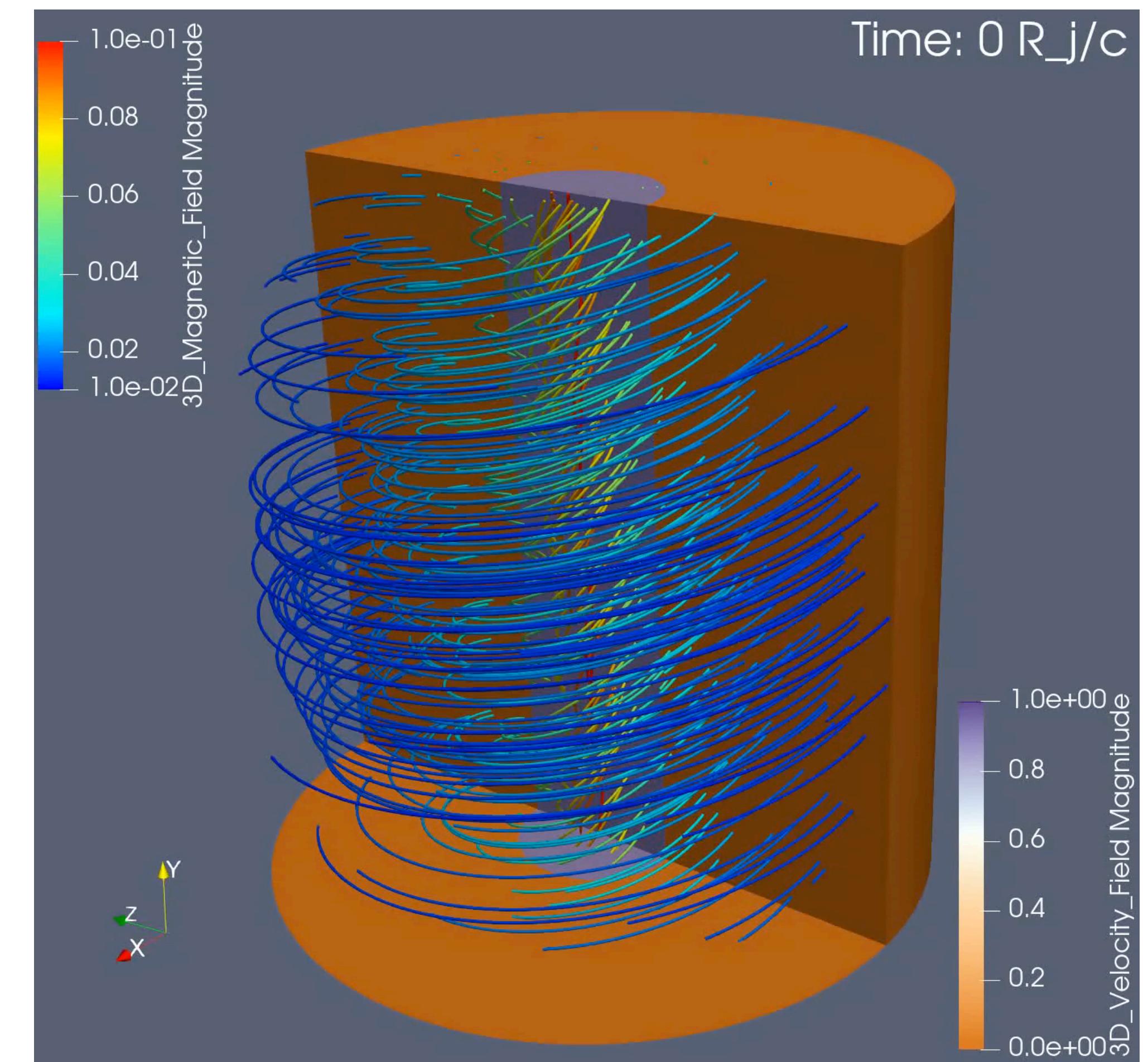
Numerical simulations

RMHD simulations with PLUTO

- Shear acceleration depends on the velocity profile, Fermi II acceleration depends on the turbulence spectrum
- Jet injected along Y axis, ambient at rest
- Periodic box along the jet axis to study the Kelvin-Helmholtz instability
- Different parameters explored with helical field
 $v \in [0.6c, 0.99c]$ & $\sigma \in [0.002, 0.2]$ $\sigma_{y,\phi} = \langle B_{y,\phi}^2 \rangle / 8\pi\rho_0 c^2$

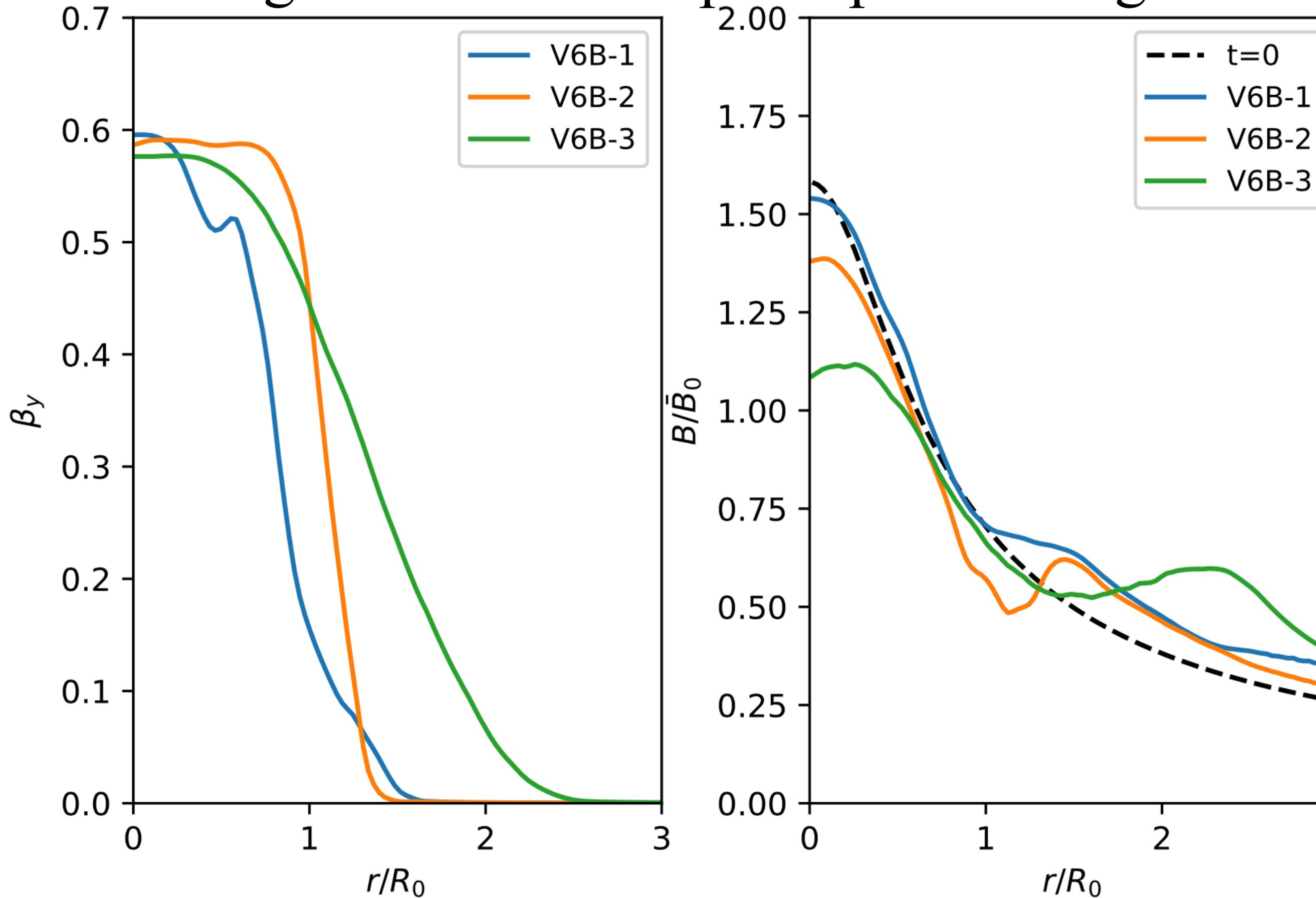
Runs*	β_0	σ_y	σ_ϕ	Box size	Grid points	Θ_0	R_0	$L_K(\text{erg s}^{-1})$
V6B-1	0.6	10^{-1}	10^{-1}	$6.0R_0$	375^3	0.01	0.1kpc	1.3×10^{43}
V6B-1-SB	0.6	10^{-1}	10^{-1}	$4.8R_0$	300^3	0.01	0.1kpc	1.3×10^{43}
V6B-1-LR	0.6	10^{-1}	10^{-1}	$6.0R_0$	200^3	0.01	0.1kpc	1.3×10^{43}
V6B-2	0.6	10^{-2}	10^{-2}	$6.0R_0$	375^3	0.01	0.1kpc	1.3×10^{43}
V6BA-2	0.6	0.016	0.004	$6.0R_0$	375^3	0.01	0.1kpc	1.3×10^{43}
V6BT-2	0.6	0.004	0.016	$6.0R_0$	375^3	0.09	0.1kpc	1.6×10^{43}
V6B-3	0.6	10^{-3}	10^{-3}	$6.0R_0$	375^3	0.01	0.1kpc	1.3×10^{43}
V9B-1	0.9	10^{-1}	10^{-1}	$8.0R_0$	500^3	0.09	1 kpc	6.7×10^{45}
V9B-2	0.9	10^{-2}	10^{-2}	$8.0R_0$	500^3	0.04	1 kpc	7.0×10^{45}
V9B-3	0.9	10^{-3}	10^{-3}	$8.0R_0$	500^3	0.02	1 kpc	6.7×10^{45}
V99B-2	0.99	10^{-2}	10^{-2}	$8.0R_0$	500^3	0.07	1 kpc	7.9×10^{46}

A sheath structure and turbulence can be self-generated via KH instability

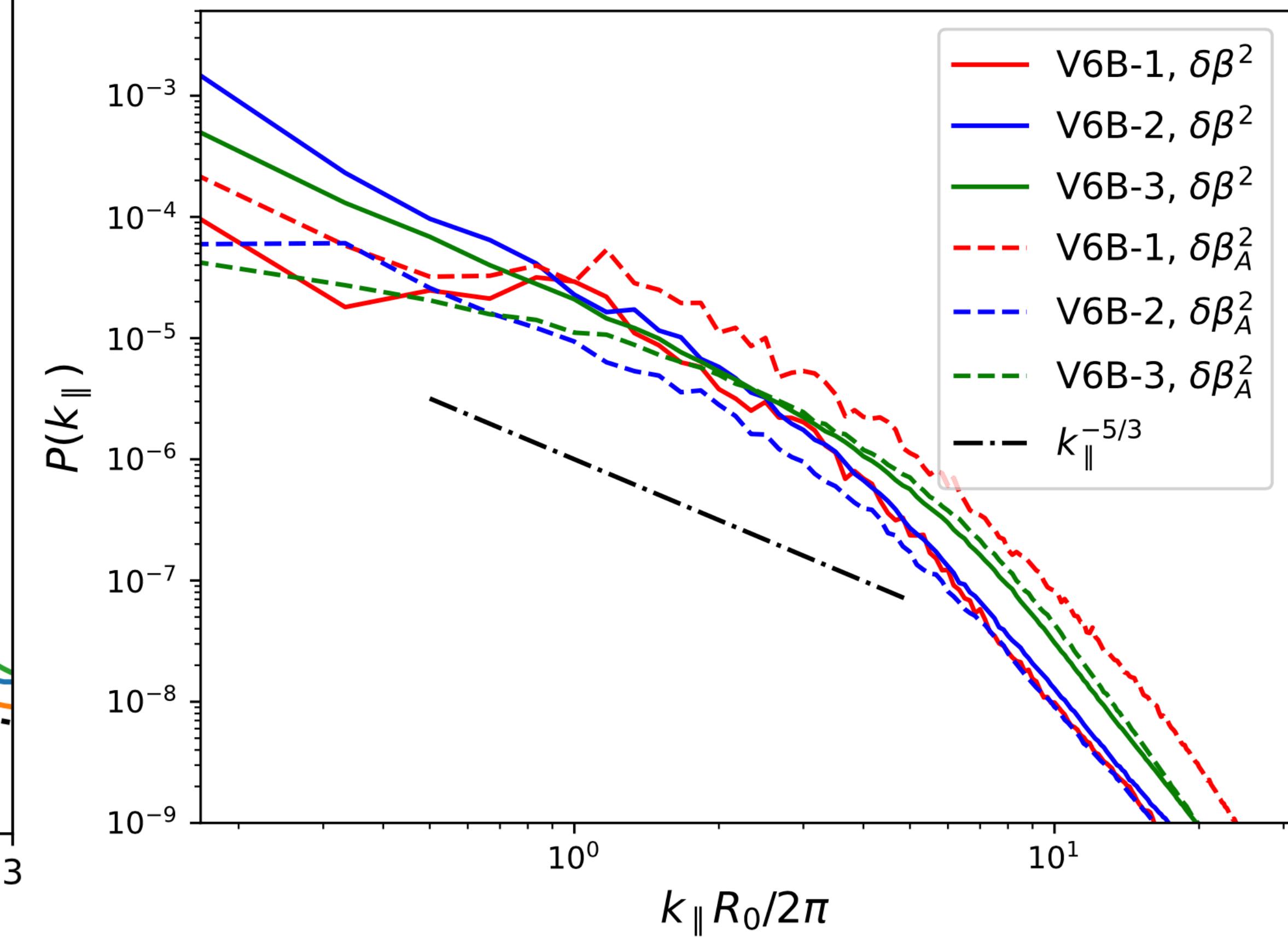


$v=0.6c$ cases in saturated KHI stage

lower magnetization case: wider sheaths
& significant B filed pile-up at the edge

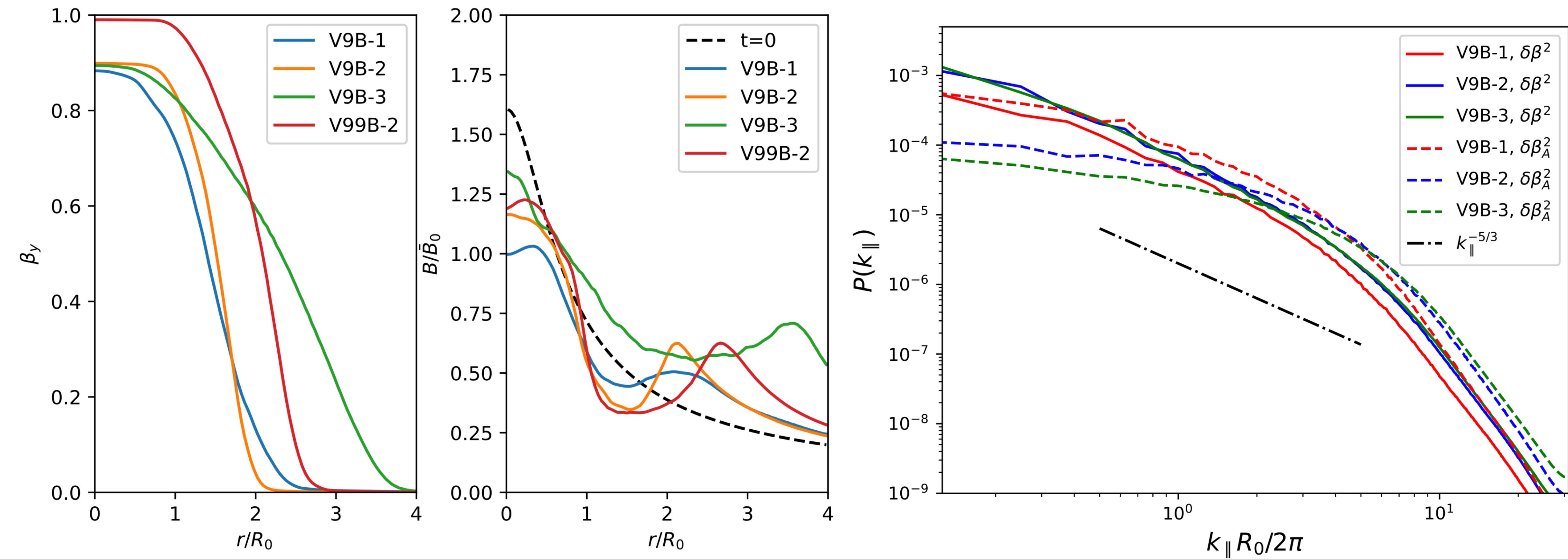


Velocity turbulence: Kolmogorov theory



$v=0.9c$ cases in saturated KHI stage

Higher velocities/lower magnetization lead to wider sheaths



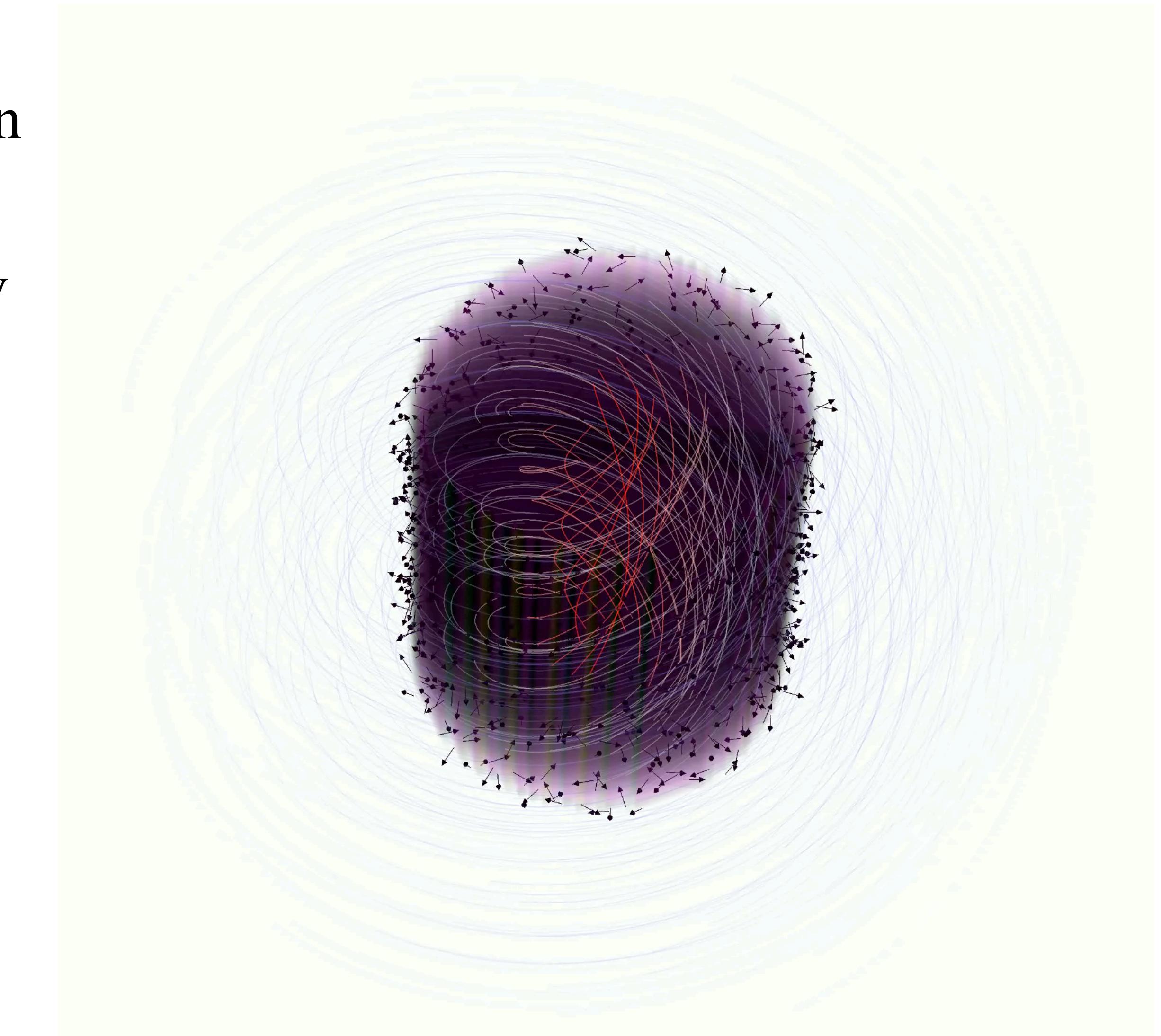
RMHD + Test-particle Simulations

- More self-consistent particle acceleration
- Higher-resolution RMHD simulation
- Inject protons with Larmor radii at a few grid scales to avoid sub-grid physics

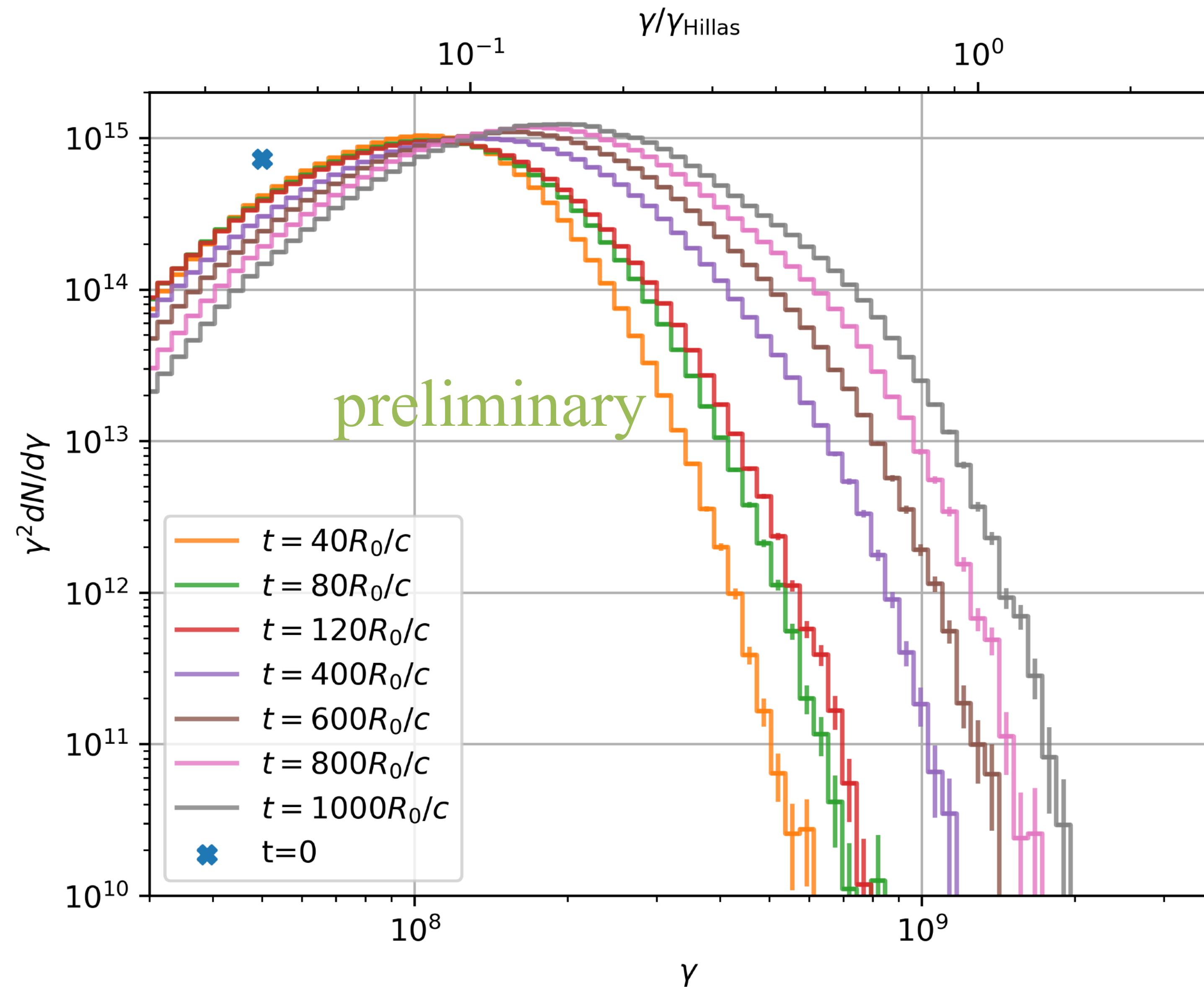
$$\frac{d\mathbf{x}_p}{dt} = \mathbf{v}_p$$

$$\frac{d(\gamma\mathbf{v})_p}{dt} = \alpha_p(c\mathbf{E} + \mathbf{v}_p \times \mathbf{B})$$

- To study the capability to accelerate UHECRs via shear acceleration



Numerical results



Summary

- **Shear (and stochastic)** acceleration is **unavoidable** in jets for **all scales**:
 - Self-generation of spine-sheath structure and turbulence via KH instability
- **Analytical solutions:** modeling of multi-wavelength SED
 - Shear acceleration can explain kpc-scale X-ray jet
 - Shear and stochastic acceleration can explain sub-pc-scale radio jet
- **Numerical studies** validates these mechanisms
 - Protons can achieve Hillas limit in jets via shear acceleration
 - Contribution to >EeV CRs from kpc-scale jets (Cen A)